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# GLOBAL EXISTENCE FOR ENERGY CRITICAL WAVES IN 3-D DOMAINS

by

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**Abstract.** — We prove that the defocusing quintic wave equation, with Dirichlet boundary conditions, is globally well posed on  $H_0^1(\Omega) \times L^2(\Omega)$  for any smooth (compact) domain  $\Omega \subset \mathbb{R}^3$ . The main ingredient in the proof is an  $L^5$  spectral projector estimate, obtained recently by Smith and Sogge [12], combined with a precise study of the boundary value problem.

## 1. Introduction

Let  $\Omega \in \mathbb{R}^3$  be a smooth bounded domain with boundary  $\partial\Omega$  and  $\Delta_D$  the Laplacian acting on functions with Dirichlet boundary conditions. We are interested in describing the relationship between certain  $L^p$  estimates for the associated spectral projector, obtained recently by H. Smith and C. Sogge [12], and Strichartz inequalities for solutions to the wave equation in  $\Omega$ . This relationship turns out to be very simple, natural and optimal (at least in some range of indexes), and is closely related to an earlier remark of Mockenhaupt, Seeger and Sogge regarding Fourier integral operators ([7], Corollary 3.3). As an application we consider the critical semi-linear wave equation (with real initial data) in  $\Omega$ ,

$$(1.1) \quad \begin{aligned} (\partial_t^2 - \Delta)u + u^5 &= 0, & \text{in } \mathbb{R}_t \times \Omega \\ u|_{t=0} &= u_0, \quad \partial_t u|_{t=0} = u_1, & u|_{\mathbb{R}_t \times \partial\Omega} = 0, \end{aligned}$$

which enjoys the conservation of energy

$$E(u)(t) = \int_{\Omega} \left( \frac{|\nabla u|^2(t, x) + |\partial_t u|^2(t, x)}{2} + \frac{|u|^6(t, x)}{6} \right) dx = E(u)(0).$$

Our main result reads:

**Theorem 1.** — *For any  $(u_0, u_1) \in H_0^1(\Omega) \times L^2(\Omega)$  there exists a unique (global in time) solution  $u$  to (1.1) in the space*

$$X = C^0(\mathbb{R}_t; H_0^1(\Omega)) \cap C^1(\mathbb{R}_t; L^2(\Omega)) \cap L_{loc}^5(\mathbb{R}; L^{10}(\Omega)).$$

**Remark 1.** — To our knowledge, the fact that weakened dispersion estimates can still imply optimal (and scale invariant) Strichartz estimates for the solution to the wave equation was first noticed by the second author [6]. Observe that the results obtained in [6], though restricted only to the interior of strictly convex domains, are far more precise than the results presented here, and apply to the critical non linear wave equation in higher dimension as well.

**Remark 2.** — The difficulty in proving Theorem 1 is that we cannot afford any loss in the Strichartz estimates we prove: Strichartz estimates obtained by Tataru [15] for Lipschitz metrics (see also the results by Anton [1] in the Schrödinger context) would certainly improve the classical result (well posedness for the cubic non linear wave equation) but it would not be enough to deal with the quintic nonlinearity.

**Remark 3.** — In other dimensions, we can still apply Smith and Sogge's spectral projectors results, leading to other Strichartz type estimates. The strategy also works for obtaining semi-classical Strichartz-type estimates for the Schrödinger equation (see Burq, Gérard and Tzvetkov [2]). These questions will be addressed elsewhere.

**Remark 4.** — Using the material in this paper, it is rather standard to prove existence of global smooth solutions, for smooth initial data satisfying compatibility conditions (see [11]). Furthermore, the arguments developed in this paper apply equally well to more general defocusing non linearities  $f(u) = V'(u)$  satisfying

$$|f(u)| \leq C(1 + |u|^5), \quad |f'(u)| \leq C(1 + |u|)^4.$$

Finally, let us remark that our results can be localized (in space) and consequently hold also in the exterior of any obstacle, and we extend in this framework previous results obtained by Smith and Sogge [11] for convex obstacles.

We shall denote in the remaining of this paper, for  $s \geq 0$ , by  $H_D^s(\Omega)$  the domain of  $(-\Delta_D)^{s/2}$  ( $H_D^s = H_0^s(\Omega)$  for  $0 \leq s < 3/2$ ).

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## 2. Local existence

The local (in time) existence result for (1.1) is in fact an easy consequence of some recent work by Smith and Sogge [12] on the spectral projector defined by  $\Pi_\lambda = 1_{\sqrt{-\Delta_D} \in [\lambda, \lambda+1]}$ .

**Theorem A (Smith-Sogge [12, Theorem 7.1]).** — *Let  $\Omega \in \mathbb{R}^3$  be a smooth bounded domain, then*

$$(2.1) \quad \|\Pi_\lambda u\|_{L^5(\Omega)} \leq \lambda^{\frac{2}{5}} \|u\|_{L^2(\Omega)}.$$

We now derive from this result some Strichartz estimates, which are optimal w.r.t scaling.

**Theorem 2.** — *Assume that for some  $2 \leq q < +\infty$ , the spectral projector  $\Pi_\lambda$  satisfies*

$$(2.2) \quad \|\Pi_\lambda u\|_{L^q(\Omega)} \leq \lambda^\delta \|u\|_{L^2(\Omega)}.$$

Then the solution to the wave equation  $v(t, x) = e^{it\sqrt{-\Delta_D}}u_0$  satisfies

$$\|v\|_{L^q((0, 2\pi)_t \times \Omega_x)} \leq C \|u_0\|_{H_D^{\delta + \frac{1}{2} - \frac{1}{q}}}.$$

*Proof.* — The proof is rather simple. In fact, one can recast the spectral projector estimate as a square function (in time) estimate for the wave equation, and use Sobolev in time, as was already observed in [7] in the context of variable coefficients wave equations (note that such square function estimates are useful in their own right in a nonlinear context: see e.g. [5], Appendix B, [14]).

Let  $(e_\lambda(x))$  be the eigenbasis of  $L^2(\Omega)$  consisting in eigenfunctions of  $-\Delta_D$  associated to the eigenvalues  $(\lambda^2)$ . Let us define on this basis an abstract self adjoint operator

$$A(e_\lambda) = [\lambda]e_\lambda$$

(where  $[\lambda]$  is the integer part of  $\lambda$ ). Now we prove the estimate with  $v$  replaced by  $\tilde{v} = e^{itA}u_0$ . We decompose  $\tilde{v}(t, x) = \sum_{k \in \mathbb{N}} v_k(t, x)$  with

$$\tilde{v}_k(t, x) = \sum_{\lambda \in \sigma(\sqrt{-\Delta_D}) \cap [k, k+1)} e^{itk} u_\lambda e_\lambda(x), \quad u_0 = \sum_{\lambda} u_\lambda e_\lambda(x).$$

Using Plancherel formula (for  $x$  fixed),

$$\|\tilde{v}(\cdot, x)\|_{H^s(0, 2\pi)}^2 = \sum_{k \in \mathbb{N}} (1+k)^{2s} \|\tilde{v}_k(\cdot, x)\|_{L^2(0, 2\pi)}^2;$$

and consequently, using Sobolev injection in the time variable for the first inequality, (2.2) in the last inequality, with  $s = \frac{1}{2} - \frac{1}{q}$ , and the fact that  $q \geq 2$  from line 3 to line 4,

$$\begin{aligned} \|\tilde{v}\|_{L^q(\Omega; L^q(0, 2\pi))}^2 &\leq C \|\tilde{v}\|_{L^q(\Omega; H^s(0, 2\pi))}^2 = \|\|\tilde{v}(\cdot, x)\|_{H^s(0, 2\pi)}\|_{L^{q/2}(\Omega)}^2 \\ &\leq C \left\| \sum_k (1+k)^{1-\frac{2}{q}} \|\tilde{v}_k(\cdot, x)\|_{L^2(0, 2\pi)} \right\|_{L^{q/2}(\Omega)}^2 \\ &\leq C \sum_k (1+|k|)^{1-\frac{2}{q}} \|\tilde{v}_k\|_{L^q(\Omega; L_t^2(0, 2\pi))}^2 \\ &\leq C \sum_k (1+|k|)^{1-\frac{2}{q}} \|\tilde{v}_k\|_{L^2((0, 2\pi); L^q(\Omega))}^2 \\ &\leq C \sum_k \sum_{\lambda \in \sigma(\sqrt{-\Delta_D}) \cap [k, k+1)} (1+|k|)^{2\delta+1-\frac{2}{q}} |u_\lambda|^2 \sim \|u_0\|_{H_D^{\delta+\frac{1}{2}-\frac{1}{q}}(\Omega)}^2. \end{aligned}$$

Consequently,  $e^{itA}$  is continuous from  $H_D^{\delta+\frac{1}{2}-\frac{1}{q}}(\Omega)$  to  $L^q((0, 2\pi); L^q(\Omega))$ . Coming back to  $v = e^{it\sqrt{-\Delta_D}}u_0$ , it satisfies

$$(i\partial_t + A)v = (A - \sqrt{-\Delta_D})v, \quad v|_{t=0} = u_0,$$

therefore, using Duhamel formula, Minkovski inequality and that both  $e^{it\sqrt{-\Delta_D}}$  and  $(A - \sqrt{-\Delta_D})$  are bounded on  $H_D^{\delta+\frac{1}{2}-\frac{1}{q}}(\Omega)$ ,

$$\begin{aligned} \|v\|_{L^q(\Omega; L^q(0, 2\pi))} &\leq C\|u_0\|_{H_D^{\delta+\frac{1}{2}-\frac{1}{q}}(\Omega)} + C\|(A - \sqrt{-\Delta_D})v\|_{L^1((0, 2\pi); H_D^{\delta+\frac{1}{2}-\frac{1}{q}}(\Omega))} \\ &\leq C'\|u_0\|_{H_D^{\delta+\frac{1}{2}-\frac{1}{q}}(\Omega)}. \end{aligned}$$

□

**Corollary 2.1.** — Consider  $u$  solution to

$$(\partial_t^2 - \Delta)u = 0, \quad u|_{\partial\Omega} = 0, \quad u|_{t=0} = u_0, \quad \partial_t u|_{t=0} = u_1.$$

Then

$$(2.3) \quad \|u\|_{L^5((0,1); L^{10}(\Omega))} \leq C (\|u_0\|_{H^1(\Omega)} + \|u_1\|_{L^2(\Omega)}).$$

As a consequence, for any initial data  $(u_0, u_1) \in H_0^1(\Omega) \times L^2(\Omega)$ , the critical non linear wave equation (1.1) is locally well posed in

$$X_T = C^0([0, T]; H_0^1(\Omega)) \cap L^5((0, T); L^{10}(\Omega)) \times C^0([0, T]; L^2(\Omega))$$

(globally for small norm initial data).

**Remark 5.** — This corollary proves the uniqueness part in Theorem 1

To prove Corollary 2.1, we observe that according to Theorems A and 2, the operator  $\mathcal{T} = e^{\pm it\sqrt{-\Delta_D}}$  satisfies

$$(2.4) \quad \|\mathcal{T}u_0\|_{L^5((0,1)\times\Omega)} \leq C\|u_0\|_{H_D^{\frac{7}{10}}(\Omega)}.$$

Applying the previous inequality to  $\Delta u_0$ , and using the  $L^p$  elliptic regularity result

$$(2.5) \quad \begin{aligned} -\Delta u + u &= f \in L^p(\Omega), \quad u|_{\partial\Omega} = 0 \Rightarrow u \in W^{2,p}(\Omega) \cap W_0^{1,p}(\Omega) \\ \text{and } \|u\|_{W^{2,p}(\Omega)} &\leq C\|f\|_{L^p(\Omega)}, \quad 1 < p < +\infty \end{aligned}$$

we get

$$(2.6) \quad \|\mathcal{T}u_0\|_{L^5((0,1); W^{2,5}(\Omega) \cap W_0^{1,5}(\Omega))} \leq C\|u_0\|_{H_D^{\frac{27}{10}}(\Omega)}$$

and consequently by (complex) interpolation between (2.4) and (2.6),

$$(2.7) \quad \|\mathcal{T}u_0\|_{L^5((0,1); W_0^{\frac{3}{10},5}(\Omega))} \leq C\|u_0\|_{H_D^1(\Omega)};$$

finally, by Sobolev embedding

$$(2.8) \quad \|\mathcal{T}u_0\|_{L^5((0,1); L^{10}(\Omega))} \leq C\|u_0\|_{H_D^1(\Omega)}.$$

To conclude, we simply observe that

$$u = \cos(t\sqrt{-\Delta_D})u_0 + \frac{\sin(t\sqrt{-\Delta_D})}{\sqrt{-\Delta_D}}u_1$$

and  $1/\sqrt{-\Delta_D}$  is an isometry from  $L^2(\Omega)$  to  $H_0^1(\Omega)$ , leading to (2.3). The remaining of Corollary 2.1 follows by a standard fixed point argument with  $(u, \partial_t u)$  in the space  $X_T$  with a sufficiently small  $T$  (depending on the initial data  $(u_0, u_1)$ ). Note that this local in time result holds irrespective of the sign of the nonlinearity.

Finally, to obtain the global well posedness result for small initial data, it is enough to remark that if the norm of the initial data is small enough, then the fixed point can be performed in  $X_{T=1}$ . Then the control of the  $H^1$  norm by the energy (which is conserved along the evolution) allows to iterate this argument indefinitely leading to global existence. Note that this result holds also irrespective of the sign of the nonlinearity because for small  $H^1$  norms, the energy always control the  $H^1$  norm.

### 3. Global existence

It turns out that our Strichartz estimates are strong enough to extend local to *global* existence for arbitrary (finite energy) data, when combined with a trace estimate and non concentration arguments.

Before going into details, let us sketch the proof. We firstly need to refine the  $L_t^5; L_x^{10}$  estimate obtained above. We shall use (2.7) instead.

The usual  $TT^*$  argument and Christ-Kiselev Lemma [3] proves the following:

**Proposition 3.1.** — *If  $u, f$  satisfy*

$$(\partial_t^2 - \Delta)u = f, \quad u|_{\partial\Omega} = 0, \quad u|_{t=0} = u_0, \quad \partial_t u|_{t=0} = u_1$$

*then*

$$(3.1) \quad \|u\|_{L^5((0,1); W_0^{\frac{3}{10},5}(\Omega))} + \|u\|_{C^0((0,1); H_0^1(\Omega))} + \|\partial_t u\|_{C^0((0,1); L^2(\Omega))} \\ \leq C \left( \|u_0\|_{H_0^1(\Omega)} + \|u_1\|_{L^2(\Omega)} + \|f\|_{L^{\frac{5}{4}}((0,1); W^{\frac{7}{10},\frac{5}{4}}(\Omega))} \right).$$

Furthermore, (3.1) holds (with the same constant  $C$ ) if one replaces the time interval  $(0,1)$  by any interval of length smaller than 1.

**Remark 6.** — An immediate consequence of (3.1) with  $f = 0$  (and Minkovski inequality) is

$$(3.2) \quad \|u\|_{L^5((0,1); W_0^{\frac{3}{10},5}(\Omega))} + \|u\|_{L^\infty((0,1); H_0^1(\Omega))} + \|\partial_t u\|_{L^\infty((0,1); L^2(\Omega))} \\ \leq C \left( \|u_0\|_{H^1(\Omega)} + \|u_1\|_{L^2(\Omega)} + \|f\|_{L^1((0,1); L^2(\Omega))} \right),$$

but we will need the  $L_t^p$  with  $p > 1$  on the righthandside of (3.1) later on.

*Proof.* — We have

$$u(t, \cdot) = \cos(t\sqrt{-\Delta_D})u_0 + \frac{\sin(t\sqrt{-\Delta_D})}{\sqrt{-\Delta_D}}u_1 + \int_0^t \frac{\sin((t-s)\sqrt{-\Delta_D})}{\sqrt{-\Delta_D}}f(s, \cdot)ds.$$

The contributions of  $(u_0, u_1)$  are easily dealt with, as previously. Let us focus on the contribution of

$$\int_0^t \frac{e^{i(t-s)\sqrt{-\Delta_D}}}{\sqrt{-\Delta_D}}f(s, \cdot)ds.$$

Denote by  $\mathcal{T} = e^{it\sqrt{-\Delta_D}}$ ; interpolating between (2.4) and (2.6),

$$\|\mathcal{T}f\|_{L^5((0,1); W^{2-\frac{7}{10},5}(\Omega) \cap W_0^{1,5}(\Omega))} \leq C \|f\|_{H_D^2(\Omega)}.$$

Let  $u_0 \in L^2$ . Then there exist  $v_0 \in H_D^2(\Omega)$  such that

$$-\Delta v_0 = u_0, \quad \|u_0\|_{L^2} \sim \|v_0\|_{H^2};$$

as a consequence, from  $\mathcal{T}u_0 = \Delta \mathcal{T}v_0$ ,

$$\begin{aligned} \|\mathcal{T}u_0\|_{L^5((0,1);W^{-\frac{7}{10},5}(\Omega))} &\leq C\|\mathcal{T}v_0\|_{L^5((0,1);W^{2-\frac{7}{10},5}(\Omega)) \cap W_0^{1,5}(\Omega)} \\ &\leq C\|v_0\|_{H_D^2(\Omega)} \sim \|u_0\|_{L^2(\Omega)}. \end{aligned}$$

By duality we deduce that the operator  $\mathcal{T}^*$  defined by

$$\mathcal{T}^*f = \int_0^1 e^{-is\sqrt{-\Delta}} f(s, \cdot) ds$$

is bounded from  $L^{\frac{5}{4}}((0,1);W^{\frac{7}{10},\frac{5}{4}}(\Omega))$  to  $L^2(\Omega)$  (observe that  $W^{\frac{7}{10},\frac{5}{4}}(\Omega) = W_0^{\frac{7}{10},\frac{5}{4}}(\Omega)$ ); using (2.7) and boundedness of  $\sqrt{-\Delta_D}^{-1}$  from  $L^2$  to  $H_D^1$ , we obtain

$$\|\mathcal{T}(\sqrt{-\Delta_D})^{-1}\mathcal{T}^*f\|_{L^5((0,1);W_0^{\frac{3}{10},5}(\Omega)) \cap L^\infty((0,1);H_0^1(\Omega))} \leq C\|f\|_{L^{\frac{5}{4}}((0,1);W_0^{\frac{7}{10},\frac{5}{4}}(\Omega))},$$

and

$$\|\partial_t \mathcal{T}(\sqrt{-\Delta_D})^{-1}\mathcal{T}^*f\|_{L^\infty((0,1);L^2(\Omega))} \leq C\|f\|_{L^{\frac{5}{4}}((0,1);W_0^{\frac{7}{10},\frac{5}{4}}(\Omega))}.$$

But

$$\mathcal{T}(\sqrt{-\Delta_D})^{-1}\mathcal{T}^*f(s, \cdot) = \int_0^1 \frac{e^{i(t-s)\sqrt{-\Delta_D}}}{\sqrt{-\Delta_D}} f(s, \cdot) ds$$

and an application of Christ-Kiselev lemma [3] allows to transfer this property to the operator

$$f \mapsto \int_0^t \frac{e^{i(t-s)\sqrt{-\Delta_D}}}{\sqrt{-\Delta_D}} f(s, \cdot) ds.$$

□

Now we remark that if  $f = u^5$ , we can estimate

$$\begin{aligned} \|u^5\|_{L^{\frac{5}{4}}((0,1);L^{\frac{30}{17}}(\Omega))} &\leq \|u\|_{L^5((0,1);L^{10}(\Omega))}^4 \|u\|_{L^\infty((0,1);L^6(\Omega))}, \\ (3.3) \quad \|\nabla_x(u^5)\|_{L^{\frac{5}{4}}((0,1);L^{\frac{10}{9}}(\Omega))} &= 5\|u^4 \nabla_x u\|_{L^{\frac{5}{4}}((0,1);L^{\frac{10}{9}}(\Omega))} \\ &\leq 5\|u\|_{L^5((0,1);L^{10}(\Omega))}^4 \|u\|_{L^\infty((0,1);H^1(\Omega))}. \end{aligned}$$

Interpolating between these two inequalities yields

$$(3.4) \quad \|u^5\|_{L^{\frac{5}{4}}((0,1);W^{\frac{7}{10},\frac{5}{4}}(\Omega))} \leq C\|u\|_{L^5((0,1);L^{10}(\Omega))}^4 \|u\|_{L^\infty((0,1);L^6(\Omega))}^{\frac{3}{10}} \|u\|_{L^\infty((0,1);H^1(\Omega))}^{\frac{7}{10}}.$$

Following ideas of Struwe [13], Grillakis [4] and Shatah-Struwe [9, 10], we will localize these estimates on small light cones and use the fact that the  $L_t^\infty; L_x^6$  norm is small in such small cones.

**Remark 7.** — In the argument above, we need (3.1) whereas (3.2) would not be sufficient.

**3.1. The  $L^6$  estimate.** — In this section we shall always consider solutions in

$$(3.5) \quad X_{<t_0} = C^0([0, t_0]; H_0^1(\Omega)) \cap L_{loc}^5([0, t_0]; L^{10}(\Omega)) \times C^0([0, t_0]; L^2(\Omega))$$

of (1.1) having bounded energy and obtained as limits in this space of smooth solutions to the analog of (1.1) where the non linearity and the initial data have been smoothed out. Consequently all the integrations by parts we will perform will be licit by a limiting argument.

*3.1.1. A priori estimate for the normal derivative.* — We start with an a priori estimate on finite energy solutions of (1.1), which is a consequence of the uniform Lopatinski condition.

**Proposition 3.2.** — *Assume that  $u$  is a weak solution to (1.1). Then we have*

$$(3.6) \quad \left\| \frac{\partial u}{\partial \nu} \right\|_{L^2((0, t_0) \times \partial \Omega)} \leq CE(u)^{1/2}$$

where  $\frac{\partial u}{\partial \nu}$  is the trace to the boundary of the exterior normal derivative of  $u$ .

*Proof.* — Take  $Z \in C^\infty(\Omega; T\Omega)$  a vector field whose restriction to  $\partial \Omega$  is equal to  $\frac{\partial}{\partial \nu}$  and compute for  $0 < T < t_0$

$$\begin{aligned} \int_0^T \int_\Omega [(\partial_t^2 - \Delta), Z]u(t, x) \cdot u(t, x) dx dt \\ = \int_0^T \int_\Omega ((\partial_t^2 - \Delta)Zu(t, x) - Z(\partial_t^2 - \Delta)u(t, x)) u(t, x) dx dt. \end{aligned}$$

Integrating by parts, we obtain

$$\begin{aligned} (3.7) \quad \int_0^T \int_\Omega [(\partial_t^2 - \Delta), Z]u(t, x) \cdot u(t, x) dx dt &= \int_0^T \int_{\partial \Omega} \left| \frac{\partial u}{\partial \nu} \right|^2 d\sigma dt \\ &+ \int_0^T \int_\Omega -(Zu)u^5(t, x) + Z(u^5)u(t, x) dx dt + \left[ \int_\Omega \partial_t(Zu) \cdot u dx \right]_0^T - \left[ \int_\Omega (Zu) \cdot \partial_t u dx \right]_0^T. \end{aligned}$$

Remark now that if  $Z = \sum_j a_j(x) \frac{\partial}{\partial x_j}$ , then integration by parts yields (using the Dirichlet boundary condition)

$$\begin{aligned} (3.8) \quad \left| \int_\Omega -(Zu)u^5(t, x) + Z(u^5)u(t, x) dx \right| &= \frac{4}{6} \left| \int_0^T \int_\Omega Z(u^6)(t, x) dx \right| \\ &= \frac{4}{6} \left| \int_0^T \int_\Omega \sum_j \frac{\partial a_j}{\partial x_j} u^6 dx \right| \leq CE(u) \end{aligned}$$

while

$$(3.9) \quad \left| \left[ \int_\Omega \partial_t(Zu)u dx \right]_0^T - \left[ \int_\Omega (Zu) \partial_t u dx \right]_0^T \right| \leq CE(u),$$

and  $[(\partial_t^2 - \Delta), Z] = -[\Delta, Z]$  as a second order differential operator in the  $x$  variable is continuous from  $H_0^1(\Omega)$  to  $H^{-1}(\Omega)$  and consequently

$$(3.10) \quad \left| \int_{t=0}^T \int_\Omega [(\partial_t^2 - \Delta), Z]u(t, x) \bar{u}(t, x) dx dt \right| \leq CE(u).$$

As the constants are uniform with respect to  $0 < T < t_0$ , collecting (3.8), (3.9) (3.10) and (3.7) yields (3.6).  $\square$

*3.1.2. The flux identity.* — By time translation, we shall assume later that  $t_0 = 0$ . Let us first define

$$\begin{aligned}
 (3.11) \quad Q &= \frac{|\partial_t u|^2 + |\nabla_x u|^2}{2} + \frac{|u|^6}{6} + \partial_t u \left( \frac{x}{t} \cdot \nabla_x \right) u, \\
 P &= \frac{x}{t} \left( \frac{|\partial_t u|^2 + |\nabla_x u|^2}{2} + \frac{|u|^6}{6} \right) + \nabla_x u \left( \partial_t u + \left( \frac{x}{t} \cdot \nabla_x \right) u + \frac{u}{t} \right), \\
 D_T &= \{x; |x| < -T\}, \\
 K_S^T &= \{(x, t); |x| < -t, S < t < T\} \cap \Omega \\
 M_S^T &= \{x; |x| = -t, S < t < T\} \\
 \partial K_S^T &= ([S, T] \times \partial\Omega) \cap K_S^T \cup D_T \cup D_S \cup M_S^T \\
 e(u) &= \left( \frac{|\partial_t u|^2 + |\nabla_x u|^2}{2} + \frac{|u|^6}{6}, -\partial_t u \nabla_x u \right)
 \end{aligned}$$

and the Flux across  $M_S^T$

$$\text{Flux } (u, M_S^T) = \int_{M_S^T} \langle e(u), \nu \rangle d\sigma(x, t)$$

where

$$\nu = \frac{1}{\sqrt{2}|x|} \left( -t, -\frac{x}{t} \right) = \frac{1}{\sqrt{2}|x|} (|x|, \frac{x}{|x|})$$

is the outward normal to  $M_S^T$  and  $d\sigma(x, t)$  the induced measure on  $M_S^T$ . Remark that

$$\begin{aligned}
 (3.12) \quad \text{Flux } (u, M_S^T) &= \int_{M_S^T} \frac{|\partial_t u|^2 + |\nabla_x u|^2}{2} + \frac{|u|^6}{6} - \partial_t u \frac{x}{|x|} \cdot \nabla_x u d\sigma(x, t) \\
 &= \int_{M_S^T} \frac{1}{2} \left| \frac{x}{|x|} \partial_t u - \nabla_x u \right|^2 + \frac{|u|^6}{6} d\sigma(x, t) \geq 0.
 \end{aligned}$$

An integration by parts gives (see Rauch [8] or [11, (3.3')])

$$\begin{aligned}
 (3.13) \quad \int_{x \in \Omega, |x| < -T} \left( \frac{|\partial_t u|^2 + |\nabla_x u|^2}{2} + \frac{|u|^6}{6} \right) (x, T) dx + \text{Flux } (u, M_S^T) \\
 = \int_{x \in \Omega, |x| < -S} \left( \frac{|\partial_t u|^2 + |\nabla_x u|^2}{2} + \frac{|u|^6}{6} \right) (x, S) dx = E_{loc}(S).
 \end{aligned}$$

This proves that  $u|_{M_S^T}$  is bounded in  $H^1(M_S^T) \cap L^6(M_S^T)$  (uniformly with respect to  $T < 0$ ) and, since  $E_{loc}(S)$  being a non-negative non-increasing function has a limit when  $S \rightarrow 0^-$ ,

$$(3.14) \quad \text{Flux}(u, M_S^0) = \lim_{T \rightarrow 0^-} \text{Flux}(u, M_S^T) = \lim_{T \rightarrow 0^-} (E_{loc}(S) - E_{loc}(T))$$

exists and satisfies

$$(3.15) \quad \lim_{S \rightarrow 0^-} \text{Flux } (u, M_S^0) = 0.$$

3.1.3. *The  $L^6$  estimate.* — We are now in position to prove the classical non concentration effect:

**Proposition 3.3.** — *Assume that  $x_0 \in \overline{\Omega}$ . Then for any solution  $u$  to (1.1) in the space  $X_{<t_0}$ , we have*

$$(3.16) \quad \lim_{t \rightarrow t_0^-} \int_{x \in \Omega \cap \{|x-x_0| < t_0-t\}} u^6(t, x) dx = 0.$$

*Proof.* — We follow [13, 4, 9, 10] and simply have to take care of the boundary terms. We can assume that  $x_0 \in \partial\Omega$  as otherwise these boundary terms disappear in the calculations below (which in this case are standard). Contrarily to [11] we cannot use any convexity assumption to obtain that these terms have the right sign, but we shall use Proposition 3.2 to control them. Performing a space-time translation, we can assume  $x_0 = 0, t_0 = 0$ . Integrating over  $K_S^T$  the identity

$$0 = \operatorname{div}_{t,x}(tQ + u\partial_t u, -tP) + \frac{|u|^6}{3},$$

we get (see [11, (3.9)–(3.12)]),

$$\begin{aligned} 0 = & \int_{D_T} (tQ + u\partial_t u)(T, x) dx - \int_{D_S} (SQ + u\partial_t u)(S, x) dx + \frac{1}{\sqrt{2}} \int_{M_S^T} (tQ + u\partial_t u + x \cdot P) d\sigma(x, t) \\ & - \int_{((S,T) \times \partial\Omega) \cap K_S^T} \nu(x) \cdot (tP) d\sigma(x, t) + \frac{1}{3} \int_{K_S^T} u^6 dx dt. \end{aligned}$$

Let  $T \rightarrow 0^-$ . Using Hölder's inequality and the conservation of energy, we get that the first term in the left tends to 0, whereas the last term is non negative. This yields

$$\begin{aligned} (3.17) \quad & - \int_{D_S} (SQ + u\partial_t u)(S, x) dx + \frac{1}{\sqrt{2}} \int_{M_S^0} (tQ + u\partial_t u + x \cdot P) d\sigma(x, t) \\ & \leq \int_{((S,0) \times \partial\Omega) \cap K_S^0} \nu(x) \cdot (tP) d\sigma(x, t). \end{aligned}$$

On the other hand, by direct calculation (see [11, (3.11)]),

$$\begin{aligned} (3.18) \quad & \frac{1}{\sqrt{2}} \int_{M_S^0} (tQ + u\partial_t u + x \cdot P) d\sigma(x, t) \\ & = \frac{1}{\sqrt{2}} \int_{M_S^0} \frac{1}{t} |t\partial_t u + x \cdot \nabla_x u + u|^2 d\sigma(x, t) + \frac{1}{2} \int_{\partial D_S} u^2(S, x) d\sigma(x) \end{aligned}$$

and (see [11, (3.12)])

$$(3.19) \quad - \int_{D_S} (SQ + u\partial_t u)(S, x) dx \geq -\frac{1}{2} \int_{\partial D_S} u^2 d\sigma(x) - S \int_{D_S} \frac{|u|^6(S, x)}{6} dx.$$

As a consequence, we obtain

$$(3.20) \quad (-S) \int_{D_S} \frac{|u|^6(S, x)}{6} dx + \frac{1}{\sqrt{2}} \int_{M_S^0} \frac{1}{t} |t \partial_t u + x \cdot \nabla_x u + u|^2 d\sigma(x, t) \\ \leq \int_{((S, 0) \times \partial\Omega) \cap K_S^0} \nu(x) \cdot t P d\sigma(x) dt$$

where  $\nu(x)$  is the exterior normal to  $\Omega$  at point  $x$  and  $d\sigma$  is the surface measure on  $\partial\Omega$ . Taking (3.11) into account (and the Dirichlet boundary condition), we obtain on  $\partial\Omega$

$$t\nu(x) \cdot P = \frac{1}{2}(\nu(x) \cdot x) \left( \frac{\partial u}{\partial \nu} \right)^2.$$

However, for  $x \in \partial\Omega$ , given that  $x_0 = 0 \in \partial\Omega$ , we have

$$\frac{x}{|x|} = t + \mathcal{O}(x), \quad \nu(x) = \nu(0) + \mathcal{O}(x)$$

where  $t$  is a unit vector tangent to  $\partial\Omega$  at  $x_0 = 0$ . Consequently, as  $\nu(0) \cdot t = 0$ ,

$$\nu(x) \cdot x = \mathcal{O}(|x|^2), \text{ for } x \in \partial\Omega$$

and the right hand side in (3.20) is bounded (using Proposition 3.2) by

$$(3.21) \quad \sup_{x \in K_S^0} |x|^2 \times \int_{(-1, 0) \times \partial\Omega} \left( \frac{\partial u}{\partial \nu} \right)^2 d\sigma(x) dt \leq C |S|^2 E(u).$$

Therefore,

$$(3.22) \quad \int_{D_S} \frac{|u|^6(S, x)}{6} dx \leq |S| E(u) + |S| \frac{1}{\sqrt{2}|S|} \int_{M_S^0} \frac{1}{|t|} |t \partial_t u + x \cdot \nabla_x u + u|^2 d\sigma(x, t);$$

finally, by Hölder's inequality and (3.12), we obtain

$$\frac{1}{\sqrt{2}|S|} \int_{M_S^0} \frac{1}{|t|} |t \partial_t u + x \cdot \nabla_x u + u|^2 d\sigma(x, t) \leq \sqrt{2} \int_{M_S^0} \frac{|x|}{|S|} \left| \frac{x}{|x|} \partial_t u - \nabla_x u \right|^2 d\sigma(x, t) \\ + \sqrt{2} \int_{M_S^0} \frac{|u|^2}{|S||t|} d\sigma(x, t) \\ \leq C \text{Flux}(u, M_S^0) + C \text{Flux}(u, M_S^0)^{1/3},$$

hence,

$$(3.23) \quad \int_{D_S} \frac{|u|^6(S, x)}{6} dx \leq |S| E(u) + C \text{Flux}(u, M_S^0) + C \text{Flux}(u, M_S^0)^{1/3}$$

for which the right hand side goes to 0 as  $S \rightarrow 0^-$  by (3.15). Remark that in the calculations above all integrals on  $K_S^0$  and  $M_S^0$  have to be understood as the limits as  $T \rightarrow 0^-$  of the respective integrals on  $K_S^T$  and  $M_S^T$  (which exist according to (3.12), (3.14)).  $\square$

**3.2. Global existence.** — In this section we consider  $u$  the unique forward maximal solution to the Cauchy problem (1.1) in the space  $X_{<t_0}$ . Assume that  $t_0 < +\infty$  and consider a point  $x_0 \in \overline{\Omega}$ ; our aim is to prove that  $u$  can be extended in a neighborhood of  $(x_0, t_0)$ , which will imply a contradiction. We perform a space time translation and assume that  $(x_0, t_0) = (0, 0)$ .

*3.2.1. Localizing space-time estimates.* — For  $t < t' \leq 0$ , let us denote by

$$\|u\|_{(L^p; L^q)(K_t^{t'})} = \left( \int_{s=t}^{t'} \left( \int_{\{|x| < -s\} \cap \Omega} |u|^q(s, x) dx \right)^{\frac{p}{q}} ds \right)^{\frac{1}{p}}$$

the  $L_t^p L_x^q$  norm on  $K_t^{t'}$  (with the usual modification if  $p$  or  $q$  is infinite). Our main result in this section reads

**Proposition 3.4.** — *For any  $\varepsilon > 0$ , there exists  $t < 0$  such that*

$$(3.24) \quad \|u\|_{(L^5; L^{10})(K_t^0)} < \varepsilon.$$

*Proof.* — We start with an extension result:

**Lemma 3.5.** — *For any  $x_0 \in \overline{\Omega}$  there exists  $r_0, C > 0$  such that for any  $0 < r < r_0$  and any  $v \in H_0^1(\Omega) \cap L^p(\Omega)$ , there exist a function  $\tilde{v}_r \in H_0^1(\Omega)$  (independent of the choice of  $1 \leq p \leq +\infty$ ), satisfying*

$$(3.25) \quad \begin{aligned} & (\tilde{v}_r - v) |_{\{|x-x_0| < r\} \cap \Omega} = 0, \\ & \int_{\Omega} |\nabla \tilde{v}|^2 \leq C \int_{\Omega} |\nabla v|^2, \quad \|\tilde{v}_r\|_{L^p(\Omega)} \leq C \|v\|_{L^p(\{|x-x_0| < r\})}. \end{aligned}$$

*In other words, we can extend functions in  $H_0^1 \cap L^p$  on the ball  $\{|x - x_0| < r\}$  to functions in  $H_0^1(\Omega) \cap L^p(\Omega)$  with uniform bounds with respect to (small)  $r > 0$ , for the  $H^1$  and the  $L^p$  norms respectively.*

*Furthermore, for any  $u \in L^\infty((-1, 0); H_0^1(\Omega)) \cap L_{loc}^1((-1, 0); L^p(\Omega))$ , there exist a function  $\tilde{u} \in L^\infty((-1, 0); H_0^1(\Omega)) \cap L_{loc}^1((-1, 0); L^p(\Omega))$ , satisfying (uniformly with respect to  $t$ )*

$$(3.26) \quad \begin{aligned} & (\tilde{u} - u) |_{\{|x-x_0| < -t\} \cap \Omega} = 0, \\ & \int_{\Omega} |\nabla \tilde{u}|^2(t, x) + |\partial_t \tilde{u}|^2(t, x) dx \leq C \int_{\Omega} |\nabla u|^2(t, x) + |\partial_t u|^2(t, x) \\ & \|\tilde{u}(t, \cdot)\|_{L^p(\Omega)} \leq C \|u(t, \cdot)\|_{L^p(\Omega \cap \{|x-x_0| < -t\})} \quad t\text{-a.s.} \end{aligned}$$

*Proof.* — Let us first prove the first part of the lemma. Let us first assume  $x_0 \in \Omega$ . We use the usual reflexion extension (suitably cut off). Fix a function  $\phi \in C_0^\infty(\frac{9}{10}, \frac{11}{10})$  equal to 1 near 1. Let us define  $\tilde{v}_r$  (in polar coordinates  $(\rho, \theta) \in (\mathbb{R}^+ \times \mathbb{S}^2)$  centered at  $x_0$ ) by

$$(3.27) \quad \tilde{v}_r(\rho, \theta) = \begin{cases} v_r(\rho, \theta) & \text{if } \rho < r \\ \phi(\frac{\rho}{r})v(2r - \rho, \theta) & \text{if } \rho > r. \end{cases}$$

An elementary calculation shows that (uniformly with respect to  $r$ )

$$\|\tilde{v}_r(x)\|_{L^p(\mathbb{R}^3)} \leq C \|u\|_{L^p(\{|x| < r\})}$$

and (using Hardy inequality to control  $\partial_x \left( \phi \left( \frac{\rho}{r} \right) \right) \times v(2r - \rho, \theta)$ ),

$$\int_{\Omega} |\nabla \tilde{v}|^2 \leq C \int_{\Omega} |\nabla v|^2.$$

Since  $x_0 \in \Omega$ , the function  $\tilde{v}_r$  is in  $H_0^1(\Omega)$  if  $r > 0$  is sufficiently small. If  $x_0 \in \partial\Omega$ , this is no longer true and we have to take care of the boundary condition  $v|_{\partial\Omega} = 0$ . To do so the idea is to take a foliation by hypersurfaces tangent to the boundary and to perform a reflexion extension *tangential* to this foliation.

We consider a change of variables  $\Psi$  such that (near  $x_0$ )

$$\Omega = \Psi(\{y = (y', y_3); y_3 > 0\}), \quad \Psi(0) = x_0, \quad \Psi'(0) = \text{Id}.$$

Writing  $y = (rz', rz_3)$ , and working in polar coordinates  $z' = (\rho, \theta) \in (\mathbb{R}^+ \times \mathbb{S})$ , the implicit function theorem (and the assumption  $\Psi'(0) = \text{Id}$ ) implies that there exists a smooth function  $\zeta(z_3, \theta, r)$  such that for  $z_3 \in [-\frac{1}{10}, \frac{1}{10}]$ ,

$$\|\Psi(r\rho\theta, rz_3)\| = r \Leftrightarrow \rho = \zeta(\theta, z_3, r).$$

We can now define an extension in a polar coordinate for  $y'$ , (for  $\frac{y_3}{r} = z_3 < \frac{1}{10}$ ) by

$$(3.28) \quad \tilde{v}_r(\rho, \theta, y_3) = \begin{cases} v(\rho, \theta, y_3) & \text{if } \rho < r\zeta(\theta, \frac{y_3}{r}, r) \\ \phi(\frac{\rho}{r})v(2r\zeta(\theta, \frac{y_3}{r}, r) - \rho, \theta, y_3) & \text{if } \rho > r\zeta(\theta, \frac{y_3}{r}, r). \end{cases}$$

Let  $\phi \in C_0^\infty(-\frac{1}{11}, \frac{1}{11})$  equal to 1 near 0. Then the extension we consider is

$$(1 - \phi)(\frac{y_3}{r})\tilde{v}_{r,1} + \phi(\frac{y_3}{r})\tilde{v}_{r,2}$$

where  $\tilde{v}_{r,1}$  is the extension we built in the interior case and  $\tilde{v}_{r,2}$  is the extension we just built (with the foliation).

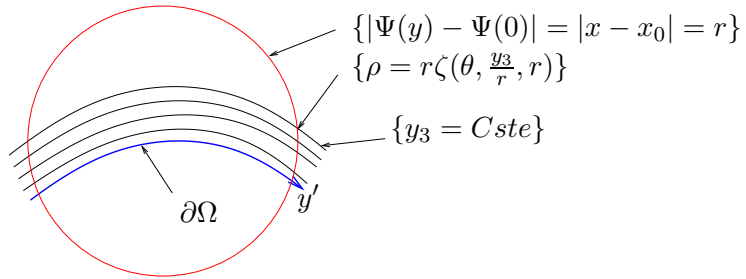


FIGURE 1. The foliation

To obtain the second part of the lemma, we just define

$$\check{u}(t, \cdot) = \tilde{u}_{-t}(t, \cdot).$$

and we can check that according to (3.27) and (3.28), it satisfies (3.26).  $\square$

Let us come back to the proof of Proposition 3.4. Let  $\check{u}$  be the function given by the second part of Lemma 3.5. Then  $(\check{u})^5$  is equal to  $u^5$  on  $K_t^0$  and

$$(3.29) \quad \begin{aligned} \|(\check{u})^5\|_{L^{\frac{5}{4}}((t,t'); L^{\frac{30}{17}}(\Omega))} &\leq \|\check{u}\|_{L^5((t,t'); L^{10}(\Omega))}^4 \|\check{u}\|_{L^\infty((t,t'); L^6(\Omega))} \\ &\leq C \|u\|_{(L^5; L^{10})(K_t^{t'})}^4 \|u\|_{(L^\infty; L^6)(K_t^0)}. \end{aligned}$$

On the other hand,  $\nabla_x(\check{u})^5 = 5(\check{u})^4 \nabla_x \check{u}$  and

$$(3.30) \quad \begin{aligned} \|\nabla_x(\check{u})^5\|_{L^{\frac{5}{4}}((t,t');L^{\frac{40}{9}}(\Omega))} &\leq 5\|\check{u}\|_{L^5((t,t');L^{10}(\Omega))}^4 \|\nabla_x \check{u}\|_{L^\infty((t,0);L^2(\Omega))} \\ &\leq C\|u\|_{(L^5;L^{10})(K_t^{t'})}^4 \|u\|_{L^\infty;H^1(\Omega)}. \end{aligned}$$

By (complex) interpolation, as in (3.4),

$$\|(\check{u})^5\|_{L^{\frac{5}{4}}((t,t');W^{\frac{7}{10},\frac{5}{4}}(\Omega))} \leq C\|u\|_{(L^5;L^{10})(K_t^{t'})}^4 \|u\|_{L^\infty;H^1(\Omega)}^{\frac{7}{10}} \|u\|_{(L^\infty;L^6)(K_t^0)}^{\frac{3}{10}}.$$

Let  $w$  be the solution (which, by finite speed of propagation, coincides with  $u$  on  $K_t^0$ ) of

$$(\partial_s^2 - \Delta)w = -(\check{u})^5, \quad w|_{\partial\Omega} = 0, \quad (w - u)|_{s=t} = \partial_s(w - u)|_{s=t} = 0,$$

applying (3.1), and the Sobolev embedding  $W^{\frac{3}{10},5}(\Omega) \hookrightarrow L^{10}(\Omega)$ , we get

$$(3.31) \quad \begin{aligned} \|u\|_{(L^5;L^{10})(K_t^{t'})} &\leq \|w\|_{L^5((t,t');L^{10}(\Omega))} \leq C\|w\|_{(L^5((t,t');W^{\frac{3}{10},5}(\Omega)))} \\ &\leq CE(u) + C\|u\|_{(L^5;L^{10})(K_t^{t'})}^4 \|u\|_{L^\infty;H^1(\Omega)}^{\frac{7}{10}} \|u\|_{(L^\infty;L^6)(K_t^0)}^{\frac{3}{10}}. \end{aligned}$$

Finally, from Proposition 3.3, (3.31) and the continuity of the mapping  $t' \in [t, 0) \rightarrow \|u\|_{(L^5;L^{10})(K_t^{t'})}$  (which takes value 0 for  $t' = t$ ), there exists  $t$  (close to 0) such that

$$\forall t < t' < 0; \|u\|_{(L^5;L^{10})(K_t^{t'})} \leq 2CE(u)$$

and passing to the limit  $t' \rightarrow 0$ ,

$$\|u\|_{(L^5;L^{10})(K_t^0)} \leq 2CE(u).$$

As a consequence, taking  $t < 0$  even smaller if necessary, we obtain

$$(3.32) \quad \|u\|_{(L^5;L^{10})(K_t^0)} \leq \varepsilon.$$

□

*3.2.2. Global existence.* — We are now ready to prove the global existence result. Let  $t < t_0 = 0$  be close to 0 and let  $v$  be the solution to the linear equation

$$(\partial_s^2 - \Delta)v = 0, \quad v|_{\partial\Omega} = 0, \quad (v - u)|_{s=t} = 0, \quad \partial_s(v - u)|_{s=t} = 0,$$

then the difference  $w = u - v$  satisfies

$$(\partial_s^2 - \Delta)w = -u^5, \quad w|_{\partial\Omega} = 0, \quad w|_{s=t} = 0, \quad \partial_s w|_{s=t} = 0.$$

Let  $\check{u}$  be the function given by Lemma 3.5 from  $u$ . We have

$$\|\check{u}\|_{L^5((t,0);L^{10}(\Omega))} \leq C\varepsilon, \quad \|\check{u}\|_{L^\infty;H^1} \leq CE(u).$$

Let  $\tilde{w}$  be the solution to

$$(\partial_s^2 - \Delta)\tilde{w} = -\check{u}^5, \quad \tilde{w}|_{\partial\Omega} = 0, \quad \tilde{w}|_{s=t} = 0, \quad \partial_s \tilde{w}|_{s=t} = 0.$$

By finite speed of propagation,  $w$  and  $\tilde{w}$  coincide in  $K_t^0$ . On the other hand, using (3.2) yields

$$(3.33) \quad \begin{aligned} \|\tilde{w}\|_{L^\infty((t,0);H^1)} + \|\partial_s \tilde{w}\|_{L^\infty((t,0);L^2(\Omega))} + \|\tilde{w}\|_{L^5((t,0);W^{\frac{3}{10},5}(\Omega))} \\ \leq C\|\check{u}^5\|_{L^1((t,0);L^2(\Omega))} \leq C\|\check{u}\|_{L^5((t,0);L^{10}(\Omega))}^5 \leq C\varepsilon^5. \end{aligned}$$

Finally, for any ball  $D$ , denote by

$$E(f(s, \cdot), D) = \int_{D \cap \Omega} (|\nabla_x f|^2 + |\partial_s f|^2 + \frac{|f|^6}{3})(s, x) dx;$$

since  $v$  is a solution to the linear equation,

$$(3.34) \quad E(v(s, \cdot), D(x_0 = 0, -s)) \rightarrow 0, \quad s \rightarrow 0^-$$

Recalling that  $u = v + \tilde{w}$  inside  $K_t^0$ , we obtain from (3.34) and (3.33) (and the Sobolev injection  $H_0^1(\Omega) \rightarrow L^6(\Omega)$ ) that there exists a small  $s < 0$  such that

$$E(u(s, \cdot), D(x_0 = 0, -s)) < \varepsilon;$$

but, since  $(u, \partial_s u)(s, \cdot) \in H_0^1(\Omega) \times L^2(\Omega)$ , we have, by dominated convergence,

$$\begin{aligned} E(u(s, \cdot), D(x_0 = 0, -s)) &= \int_{\Omega} 1_{\{|x-x_0| < -s\}}(x) (|\nabla u(s, x)|^2 + |\partial_s u(s, x)|^2 + \frac{|u|^6(s, x)}{3}) dx \\ \text{and } \lim_{\alpha \rightarrow 0} \int_{\Omega} 1_{\{|x-x_0| < \alpha-s\}}(x) (|\nabla u(s, x)|^2 + |\partial_s u(s, x)|^2 + \frac{|u|^6(s, x)}{3}) dx \\ &= \lim_{\alpha \rightarrow 0} E(u(s, \cdot), D(x_0 = 0, -s + \alpha)); \end{aligned}$$

consequently, there exists  $\alpha > 0$  such that

$$E(u(s, \cdot), D(x_0 = 0, -s + \alpha)) \leq 2\varepsilon.$$

Now, according to (3.13), the  $L^6$  norm of  $u$  remains smaller than  $2\varepsilon$  on  $\{|x - x_0| < \alpha - s'\}$ ,  $s \leq s' < 0$ . As a consequence, the same proof as for Proposition 3.4 shows that the  $L^5; L^{10}$  norm of the solution on the truncated cone

$$K = \{(x, s'); |x - x_0| < \alpha - s', s < s' < 0\}$$

is bounded. Since this is true for all  $x_0 \in \overline{\Omega}$ , a compactness argument shows that

$$\|u\|_{L^5((s,0); L^{10}(\Omega))} < +\infty$$

which, by Duhamel formula shows that

$$\lim_{s' \rightarrow 0^-} (u, \partial_s u)(s', \cdot)$$

exists in  $(H_0^1(\Omega) \times L^2(\Omega))$  and consequently  $u$  can be extended for  $s' > 0 = t_0$  small enough, using Corollary 2.1.

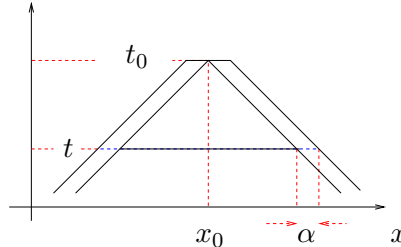


FIGURE 2. The truncated cone

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